

Necessary Optimality Conditions for Determining of the Position of the Boundary of Oil Deposit

N.K. Shazhdekeyeva, B. Kenzhegulov, A.N. Myrzasheva, G.T. Kabylkhamitov, R.U. Tuleuova

Department of Mathematics and Methods of Teaching Mathematics, Atyrau State University named after Kh.Dosmukhamedov, Student Ave.,
 1 – 060011, Atyrau, Republic of Kazakhstan.

Abstract:

The essential criteria was discovered in the thesis to define free boundary over the nonlinear partial differential equations. The task is to find co-ordinates of each point of the border so that to minimize functional. Such task can be treated as an optimum control problem. It is a problem with restriction. In accordance with a method of the interfaced functions we will pass to a problem of optimization without restrictions, having entered new functional. The purpose of the given work is calculation of the first variation of functional. It also defines a necessary condition of optimality for definition of position of border of an oil layer. The considered method for determining free boundaries is based on solving problems of determining the boundary of a system that describes a nonlinear partial differential equation.

Keywords - Differential Equation, Functional, Dirac Function, Ecessary Optimality, Initial and Boundary Conditions, Gradient, Laplacian of the Functions, Variation, Integration.

I. INTRODUCTION

The task of the current study is to find co-ordinates of each point of the border of the oil deposit so that to minimize functional. Such task can be treated as an optimum control problem. It is a problem with restriction.

The purpose of the given work is calculation of the first variation of functional.

This subsection considers the necessary optimality condition for determining of the position of the boundary of the oil deposit.

Let us consider the problem of determining of the boundary for a system described by nonlinear differential equations in partial derivatives.

$$u_t = f(x, y, t, u, u_x, u_y, u_{xx}, u_{yy}) \quad (1)$$

where $u(x, y, t)$ – m -dimensional a vector function of the state of the system, defined in a two-dimensional region $\sigma \in E^2$; u_t, u_x, u_{xx} and so on - denoting the partial derivatives with respect to time and spatial coordinates, respectively.

The initial and boundary conditions are given in general form

$$u(x, y, 0) = u_0(x, y) \quad (x, y) \in \sigma \quad (2)$$

$$g(t, x, y, u, u_n) = 0 \quad (x, y) \in \partial\sigma \quad (3)$$

where u_n denotes the normal derivative of the vector u to the boundary of the region $\partial\sigma$.

II. MATERIAL AND METHODS

In accordance with a method of the interfaced functions it is supposed to pass to a problem of optimization without restrictions, having entered new functional.

The considered method for determining free boundaries is based on solving problems of determining the boundary of a system that describes a nonlinear partial differential equation. Principles for solving ill-defined problems are considered in the works of Zakirov and Lapuk [1], Lions [2], Bubnov [3], Bulygin [4-5], Gutnikov et al. [6-8].

Applied to the inverse problems of the filtration theory in the works of Zhirov [9], Shazhdekeyeva and Mukhambetzhano [10], Shazhdekeyeva et al. [11], Frolov [12].

III. RESULTS AND DISCUSSION

The system evolves over a period of time $t \in [0, T]$, during which the measurements are made.

Observations are represented by the q -dimensional ($q < m$) vector $z(x, y, t)$, which for convenience can be considered a continuous function $z \in E^q(\sigma \times [0, T])$. Suppose that the vector z is related to the state of the system by the relation

$$z(x, y, t) = h(u) + \xi(x, y, t), \quad (4)$$

where $\xi(x, y, t)$ – measurement error. The problem is to determine the region σ (i.e., finding the coordinates of each boundary point $\partial\sigma$) in such a way as to minimize the functional

$$J = \int_0^T \int_{\sigma} [z(x, y, t) - h(u(x, y, t))]^T G(x, y, \xi, \eta, t) \times [z(\xi, \eta, t) - h(u(\xi, \eta, t))] dx dy d\xi d\eta dt \quad (5)$$

The weight matrix $G(x, y, \xi, \eta, t)$ is continuous in its arguments, positive definite and symmetric ($G = G'$). In this case, the problem under consideration can be treated as an optimal control problem, in which the position of the boundary is a control variable.

In practice, the case of measuring the state vector of a system in M-discrete points (x_j, y_j) is often encountered, i.e. a case

$$z(x_j, y_j, t) = A_j u(x_j, y_j, t) + \zeta_j(t), j = 1, 2, \dots, M \quad (6)$$

where $A_j (j = 1, 2, \dots, M)$ – matrices of dimension $(q \times m)$, whose elements are equal to zero or one. In this case, the original functional J can be represented as follows:

$$J = \int_0^T \sum_{j=1}^M [z(x_j, y_j, t) - A_j u(x_j, y_j, t)]' G_j(t) [z(x_j, y_j, t) - A_j u(x_j, y_j, t)] dt \quad (7)$$

The notation (1.5.7) can be reduced to a more general form (5) by introducing into the weight function σ - a Dirac function. Therefore, in the sequel, for convenience, a more general formula (5) is used.

The optimal control problem formulated above is a problem with the constraint imposed by equation (1.5.1). In accordance with the method of conjugate functions, pass to the optimization problem without restrictions, introducing a new functional

$$J_\psi = J + \int_0^T \int_\sigma \psi'(x, y, t) [f - u_t] dx dy dt, \quad (8)$$

where $\psi(x, y, t)$ – m-dimensional conjugate vector-valued function. If $\partial\sigma$ -boundary, minimizing J, then it also minimizes and J_ψ .

Suppose that the boundary $\partial\sigma$ is subjected to a perturbation and transforms $d\sigma^* = d\sigma + \delta(d\sigma)$ with the corresponding transformation of the region σ in $\sigma^* = \sigma + \delta\sigma$. Perturbation of the border $\partial\sigma$ leads to perturbation J_ψ . Our goal is to calculate the first variation J_ψ .

Suppose that a new area σ^* (its coordinates are indicated x^*, y^*), which depends on the parameter ε , can be converted to the original region σ by transformations

$$x^* = \Phi_1(x, y, u, \nabla u, \nabla^2 u; \varepsilon); \quad (9)$$

$$y^* = \Phi_2(x, y, u, \nabla u, \nabla^2 u; \varepsilon), \quad (10)$$

where ∇u and $\nabla^2 u$ gradient, and Laplacian of the functions u . The new value of the function $u^* = u^*(x^*, y^*)$ is reduced to the original function $u(x, y)$ by the transformations (9), (10) and the transformation

$$u^* = \Phi_3(x, y, u, \nabla u, \nabla^2 u; \varepsilon) \quad (11)$$

It is assumed that these transformations are continuous, invertible, differentiable, and that to the values $\varepsilon = 0$ correspond identical transformations

$$x = \Phi_1(x, y, u, \nabla u, \nabla^2 u; 0) \quad (12)$$

$$y = \Phi_2(x, y, u, \nabla u, \nabla^2 u; 0) \quad (13)$$

$$u = \Phi_3(x, y, u, \nabla u, \nabla^2 u; 0) \quad (14)$$

If ε - a small value, then equation (9) - (11) can be represented in the form

$$x^* = \Phi_1|_{\varepsilon=0} + \varepsilon \frac{\partial \Phi_1}{\partial \varepsilon} |_{\varepsilon=0} + O(\varepsilon) = x + \varepsilon \rho_1 + O(\varepsilon); \quad (15)$$

$$y^* = \Phi_2|_{\varepsilon=0} + \varepsilon \frac{\partial \Phi_2}{\partial \varepsilon} |_{\varepsilon=0} + O(\varepsilon) = y + \varepsilon \rho_2 + O(\varepsilon); \quad (16)$$

$$u^* = \Phi_3|_{\varepsilon=0} + \varepsilon \frac{\partial \Phi_3}{\partial \varepsilon} |_{\varepsilon=0} + O(\varepsilon) = u + \varepsilon \rho_3 + O(\varepsilon); \quad (17)$$

The first variations x, y, u are defined as follows:

$$\delta x = x^* - x = \varepsilon \rho_1; \quad (18)$$

$$\delta y = y^* - y = \varepsilon \rho_2; \quad (19)$$

$$\delta u = u^*(x^*, y^*) - u(x, y) = \varepsilon \rho_3 \quad (20)$$

The first variation J_ψ due to the perturbation is the principal linear part of (relative to ε) of the difference:

$$J_\psi[u^*(x^*, y^*)] - J_\psi[u(x, y)] = \int_0^T \int_{\sigma^*} [z - h(u^*)]' \times G(x^*, y^*, \xi^*, \eta^*, t) [z - h(u^*)] dx^* dy^* d\xi^* d\eta^* dt - \int_0^T \int_{\sigma} [z - h(u)]' G(x, y, \xi, \eta, t) [z - h(u)] \times dx dy d\xi d\eta dt + \int_0^T \int_{\sigma^*} \psi' [f - u_t^*] dx^* dy^* dt - \int_0^T \int_{\sigma} \psi' [f - u_t] dx dy dt. \quad (21)$$

Using the Jacobian of the transformation, we reduce equation (21) to the form

$$J_\psi[u^*] - J_\psi[u] = \int_0^T \int_{\sigma^*} [z - h(u^*)]' G [z - h(u^*)] \times \left| \frac{\partial(x^*, y^*, \xi^*, \eta^*)}{\partial(x, y, \xi, \eta)} \right| dx dy d\xi d\eta dt - \int_0^T \int_{\sigma} [z - h(u)]' G \times [z - h(u)] dx dy d\xi d\eta dt + \int_0^T \int_{\sigma^*} \psi' [f - u_t^*] \times \left| \frac{\partial(x^*, y^*)}{\partial(x, y)} \right| dx dy dt - \int_0^T \int_{\sigma} \psi' [f - u_t] dx dy dt \quad (22)$$

where

$$\left| \frac{\partial(x^*, y^*, \xi^*, \eta^*)}{\partial(x, y, \xi, \eta)} \right| \approx 1 + \varepsilon \frac{\partial \varphi_1}{\partial x} + \varepsilon \frac{\partial \varphi_2}{\partial y} + \varepsilon \frac{\partial \varphi_1}{\partial \xi} + \varepsilon \frac{\partial \varphi_2}{\partial \eta},$$

$$\left| \frac{\partial(x^*, y^*)}{\partial(x, z)} \right| \approx 1 + \varepsilon \frac{\partial \varphi_1}{\partial x} + \varepsilon \frac{\partial \varphi_2}{\partial y}.$$

Expanding the integrals in the expression (22) in a Taylor series and preserving terms of the first order with respect to ε , obtain the first variation J_ψ in the form

$$\begin{aligned} \delta J_\psi = & - \int_0^T \int \int \int_{\sigma} 2[z(\xi, \eta, t) - h(u(\xi, \eta, t))] G'(x, y, \xi, \eta, t) \times h_u(u(x, y, t)) [\bar{\delta u} + u_x \delta x + \\ & + u_y \delta y] dx dy d\xi d\eta dt + \int_0^T \int \int \int_{\sigma} 2[z(x, y, t) - h(u(x, y, t))] G(x, y, \xi, \eta, t) \times \\ & \times [z(\xi, \eta, t) - h(u(\xi, \eta, t))] \left[\frac{\partial \delta x}{\partial x} + \frac{\partial \delta y}{\partial y} \right] dx dy d\xi d\eta dt + \int_0^T \int \int_{\sigma} [\psi'(f - u_x) \left(\frac{\partial \delta x}{\partial x} + \frac{\partial \delta y}{\partial y} \right) + \\ & + \delta \psi'(f - u_x) + \psi' f_x \delta x + \psi' f_y \delta y + \psi' f_u \bar{\delta u} + \psi' f_{u_x} u_x \delta x + \psi' f_{u_y} u_y \delta y + \psi' f_{u_x} \bar{\delta u}_x + \\ & + \psi' f_{u_x} u_{xy} \delta y + \psi' f_{u_x} \bar{\delta u}_y + \psi' f_{u_x} u_{yx} \delta x + \psi' f_{u_x} u_{yy} \delta y + \psi' f_{u_x} \bar{\delta u}_{xx} + \psi' f_{u_x} u_{xxx} \delta x + \psi' f_{u_x} u_{xy} \delta y + \\ & + \psi' f_{u_x} u_{xy} \delta y + \psi' f_{u_x} \bar{\delta u}_{yy} + \psi' f_{u_x} u_{yxx} \delta x + \psi' f_{u_x} u_{yyy} \delta y - \psi' \bar{\delta u}_x - \psi' u_x \delta x - \psi' u_y \delta y] dx dy dt, \end{aligned} \quad (23)$$

there

$$\delta \psi = \psi(x^*, y^*, t) - \psi(x, y, t) = \psi_x \delta x + \psi_y \delta y, \quad \bar{\delta u} = u^*(x, y, t) - u(x, y, t), \quad \psi' \bar{\delta u}_t = \frac{\partial}{\partial t} (\psi' \bar{\delta u}) - \frac{\partial \psi'}{\partial t} \bar{\delta u}; \quad (27)$$

The following relations are also used:

$$\begin{aligned} \psi' f_{u_{xx}} \bar{\delta u}_{xx} + \psi' f_{u_{yy}} \bar{\delta u}_{yy} = & - \frac{\partial}{\partial x} [(\psi' f_{u_{xx}})_x \bar{\delta u}] + \frac{\partial^2}{\partial x^2} (\psi' f_{u_{xx}}) \bar{\delta u} + \frac{\partial}{\partial x} (\psi' f_{u_{xx}} \bar{\delta u}_x) - \frac{\partial}{\partial y} [(\psi' f_{u_{yy}})_y \bar{\delta u}] + \\ & + \frac{\partial^2}{\partial y^2} (\psi' f_{u_{yy}}) \bar{\delta u} + \frac{\partial}{\partial y} (\psi' f_{u_{yy}} \bar{\delta u}_y) \end{aligned} \quad (28)$$

and the relation $\bar{\delta u}(x, y, 0) = 0$, obtain the first variation of the functional δJ_ψ in the form

$$\begin{aligned} \delta J_\psi = & \int_0^T \int \int \int_{\sigma} \left\{ \frac{\partial \psi'}{\partial t} + \psi' f_u - \frac{\partial}{\partial x} (\psi' f_{u_x}) - \frac{\partial}{\partial y} (\psi' f_{u_y}) + \frac{\partial^2}{\partial x^2} (\psi' f_{u_{xx}}) + \frac{\partial^2}{\partial y^2} (\psi' f_{u_{yy}}) - 2 \int_{\sigma} [z(\xi, \eta, t) - \right. \\ & \left. - h(u(\xi, \eta, t))] G'(x, y, \xi, \eta, t) h_u(u(x, y, t)) d\xi d\eta \right\} \times \bar{\delta u} dx dy dt + \int_0^T \int \int_{\sigma} \left\{ \frac{\partial}{\partial x} [\psi' f_{u_x} \bar{\delta u} - (\psi' f_{u_x})_x \bar{\delta u} + \right. \\ & \left. + \psi' f_{u_x} \bar{\delta u}_x + \int_{\sigma} [z(x, y, t) - h(u(x, y, t))] \times G(x, y, \xi, \eta, t) [z(\xi, \eta, t) - h(u(\xi, \eta, t))] d\xi d\eta \right\} \delta x + \\ & \left. + \frac{\partial}{\partial y} [\psi' f_{u_y} \bar{\delta u} - (\psi' f_{u_y})_y \bar{\delta u} + \psi' f_{u_y} \bar{\delta u}_y + 2 \int_{\sigma} [z(x, y, t) - h(u(x, y, t))] G(x, y, \xi, \eta, t) \times [z(\xi, \eta, t) - \right. \end{aligned}$$

$$\begin{aligned} \delta u &= \bar{\delta u} + u_x \delta x + u_y \delta y; \\ \delta u_x &= \frac{\partial u^*(x^*, y^*, t)}{\partial x^*} - \frac{\partial u(x, y, t)}{\partial x} = \frac{\partial u^*(x, y, t)}{\partial x} - \frac{\partial u(x, y, t)}{\partial x} + u_{xx} \delta x + u_{xy} \delta y = \bar{\delta u}_x + u_{xx} \delta x + u_{xy} \delta y; \\ \delta u_y &= \bar{\delta u}_y + u_{yx} \delta x + u_{yy} \delta y; \\ \delta u_{xx} &= \bar{\delta u}_{xx} + u_{xxx} \delta x + u_{xxy} \delta y; \\ \delta u_{yy} &= \bar{\delta u}_{yy} + u_{yyx} \delta x + u_{yyy} \delta y. \end{aligned}$$

In addition, it is assumed that if the point (x, y) does not fall in the region σ^* , then $u^*(x, y)$ it can still be represented in the form $u^*(x^*, y^*)$, where (x^*, y^*) - is the point in σ^* , which corresponds to the point (x, y) be reason of the transformations (9), (10).

Using equations

$$\begin{aligned} \frac{\partial}{\partial x} (\psi' f \delta x) + \frac{\partial}{\partial y} (\psi' f \delta y) = & \psi' f \delta x_x + \psi' f \delta y_y + \frac{\partial \psi'}{\partial x} f \delta x + \frac{\partial \psi'}{\partial y} f \delta y + \psi' f_x \delta x + \psi' f_y \delta y + \psi' f_{u_x} u_x \delta x + \psi' f_{u_y} u_y \delta y + \\ & + \psi' f_{u_x} u_{yx} \delta x + \psi' f_{u_x} u_{xxx} \delta x + \psi' f_{u_x} u_{yxx} \delta x + \psi' f_y \delta y + \psi' f_{u_x} u_y \delta y + \psi' f_{u_x} u_{xy} \delta y + \psi' f_{u_x} u_{yy} \delta y + \\ & + \psi' f_{u_x} u_{xy} \delta y + \psi' f_{u_x} \bar{\delta u}_y + \psi' f_{u_x} u_{xy} \delta y + \psi' f_{u_x} u_{yxx} \delta x + \psi' f_{u_x} u_{yyy} \delta y; \end{aligned} \quad (24)$$

$$\frac{\partial}{\partial x} (\psi' u_x \delta x) + \frac{\partial}{\partial y} (\psi' u_y \delta y) = \frac{\partial \psi'}{\partial x} u_x \delta x + \psi' u_{xx} \delta x + \psi' u_x \delta x_x + \frac{\partial \psi'}{\partial y} u_y \delta y + \psi' u_{yy} \delta y + \psi' u_y \delta y_y; \quad (25)$$

$$\psi' f_{u_x} \bar{\delta u}_x + \psi' f_{u_x} \bar{\delta u}_y = \frac{\partial}{\partial x} (\psi' f_{u_x} \bar{\delta u}) - \frac{\partial}{\partial x} (\psi' f_{u_x}) \bar{\delta u} + \frac{\partial}{\partial y} (\psi' f_{u_x} \bar{\delta u}) - \frac{\partial}{\partial y} (\psi' f_{u_x}) \bar{\delta u}; \quad (26)$$

$$-h(u(\xi, \eta, t))d\xi d\eta]dxdydt - \int_{\sigma} \psi' \overline{\delta u} \Big|_{t=T} dx dy. \quad (29)$$

Let the conjugate vector $\psi(x, y, t)$ satisfy the system of equations

$$\psi_t = -f'_u \psi + (f'_{u_x} \psi)_x + (f'_{u_y} \psi)_y - (f'_{u_{xx}} \psi)_{xx} - (f'_{u_{yy}} \psi)_{yy} + 2 \int_{\sigma} h'_u(x, y, t) G(x, y, \xi, \eta, t) \times [z(\xi, \eta, t) - h(u(\xi, \eta, t))] d\xi d\eta \quad (30)$$

with the condition at the end of the time interval

$$\psi(x, y, T) = 0. \quad (31)$$

The substitution of the equations (30) and (31) into the expression (29) and the application of the Gauss-Ostrogradskii theorem gives the following representation for the variation of the functional:

$$\delta J_{\psi} = \int_{0}^T \int_{\partial \sigma} \{[\psi' f'_{u_x} - (\psi' f'_{u_x})_x] \cos \nu + [\psi' f'_{u_y} - (\psi' f'_{u_y})_y] \sin \nu\} \overline{\delta u} ds dt + \int_{0}^T \int_{\partial \sigma} (\psi' f'_{u_x} \overline{\delta u}_x \cos \nu + \psi' f'_{u_y} \overline{\delta u}_y \sin \nu) ds dt + \int_{0}^T \int_{\sigma} \{2[z(x, y, t) - h(u(x, y, t))] G(x, y, \xi, \eta, t) \times [z(\xi, \eta, t) - h(u(\xi, \eta, t))] d\xi d\eta\} [\delta x \cos \nu + \delta y \sin \nu] ds dt, \quad (32)$$

where ν – the angle between the positive direction of the x axis and the outer normal to the boundary of the region $\partial \sigma$, s – the arc length of the contour $\partial \sigma$.

To determine the boundary conditions for a function ψ it is more convenient to use a $\overline{\delta u}$. Expressing $\overline{\delta u}$ through δu and substituting this expression in (32), obtain

$$\delta J_{\psi} = \int_{0}^T \int_{\partial \sigma} \{[\psi' f'_{u_x} - (\psi' f'_{u_x})_x] \cos \nu + [\psi' f'_{u_y} - (\psi' f'_{u_y})_y] \sin \nu\} \delta u ds dt + \int_{0}^T \int_{\partial \sigma} (\psi' f'_{u_x} \overline{\delta u}_x \cos \nu + \psi' f'_{u_y} \overline{\delta u}_y \sin \nu) ds dt + \int_{0}^T \int_{\partial \sigma} \{S(x, y, t)(\delta x \cos \nu + \delta y \sin \nu) - [(\psi' f'_{u_x} - (\psi' f'_{u_x})_x) \cos \nu + (\psi' f'_{u_y} - (\psi' f'_{u_y})_y) \sin \nu](u_x \delta x + u_y \delta y)\} ds dt, \quad (33)$$

where is denoted

$$S(x, y, t) = 2 \int_{\sigma} [z(x, y, t) - h(u(x, y, t))] G(x, y, \xi, \eta, t) [z(\xi, \eta, t) - h(u(\xi, \eta, t))] d\xi d\eta$$

If we express both $\overline{\delta u}_x$ and $\overline{\delta u}_y$ in terms of the normal $\overline{\delta u}_n$ and tangential $\overline{\delta u}_s$ derivatives of the function u on $\partial \sigma$, and also assume that the contour $\partial \sigma$ is closed, then equation (1.5.33) is transformed to the form

$$\delta J_{\psi} = \int_{0}^T \int_{\partial \sigma} \{[\psi' f'_{u_x} - (\psi' f'_{u_x})_x] \cos \nu + [\psi' f'_{u_y} - (\psi' f'_{u_y})_y] \sin \nu +$$

$$+ \frac{\partial}{\partial x} [\psi' f'_{u_{xx}} \sin \nu \cos \nu - \psi' f'_{u_{yy}} \cos \nu \sin \nu] \delta u ds dt + \int_{0}^T \int_{\partial \sigma} \{S(x, y, t)(\delta x \cos \nu + \delta y \sin \nu) - [(\psi' f'_{u_x} - (\psi' f'_{u_x})_x) \cos \nu + (\psi' f'_{u_y} - (\psi' f'_{u_y})_y) \sin \nu][u_x \delta x + u_y \delta y] - \frac{\partial}{\partial s} (\psi' f'_{u_{xx}} \sin \nu \cos \nu - \psi' f'_{u_{yy}} \cos \nu \sin \nu)(u_x \delta x + u_y \delta y)\} ds dt \quad (34)$$

For convenience of calculations, the variation $\overline{\delta u}_n = \partial u^*(x, y) / \partial n - \partial u(x, y) / \partial n$ can be expressed in terms of $\delta u_n = \partial u^*(x^*, y^*) / \partial n^* - \partial u(x, y) / \partial n$.

Indeed, since the angle between n and n^* , obtain

$$\frac{\partial u^*(x, y)}{\partial n} - \frac{\partial u(x, y)}{\partial n} = \frac{\partial u^*(x^*, y^*)}{\partial n^*} - \frac{\partial u(x, y)}{\partial n} - \frac{\partial^2 u(x, y)}{\partial n^2} \delta n - \frac{\partial^2 u(x, y)}{\partial n \partial s} \delta s, \quad (35)$$

where n – normal direction to $\partial \delta$ at the point (x, y) ; s – the tangent direction to $\partial \delta$ at the point (x, y) ; n^* – the normal direction to the perturbed boundary $\partial \sigma^*$ at the point (x^*, y^*) , which after deformation corresponds to the point (x, y) on $\partial \delta$; ∂n and ∂s – are the normal and tangential components of the variations δx and δy .

The substitution of (35) into (34) gives

$$\delta J_{\psi} = \int_{0}^T \int_{\partial \sigma} \{[\psi' f'_{u_x} - (\psi' f'_{u_x})_x] \cos \nu + [\psi' f'_{u_y} - (\psi' f'_{u_y})_y] \sin \nu + \frac{\partial}{\partial x} [\psi' f'_{u_{xx}} \sin \nu \cos \nu - \psi' f'_{u_{yy}} \cos \nu \sin \nu] \delta u ds dt + \int_{0}^T \int_{\partial \sigma} (\psi' f'_{u_{xx}} \cos^2 \nu + \psi' f'_{u_{yy}} \sin^2 \nu) \delta u_n ds dt - \int_{0}^T \int_{\partial \sigma} (\psi' f'_{u_{xx}} \cos^2 \nu + \psi' f'_{u_{yy}} \sin^2 \nu) (\frac{\partial^2 u}{\partial n^2} \delta n + \frac{\partial^2 u}{\partial n \partial s} \delta s) ds dt + \int_{0}^T \int_{\partial \sigma} \{S(x, y, t)(\delta x \cos \nu + \delta y \sin \nu) - [(\psi' f'_{u_x} - (\psi' f'_{u_x})_x) \cos \nu + (\psi' f'_{u_y} - (\psi' f'_{u_y})_y) \sin \nu][u_x \delta x + u_y \delta y] -$$

$$-\frac{\partial}{\partial S}(\psi' f_{u_{xx}} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu)(u_x \delta x + u_y \delta y) \} ds dt \quad (36)$$

It follows from (3) that

$$g_x \delta x + g_y \delta y + g_u \delta u + g_{u_n} \delta u_n = 0, \\ x, y \in \partial \sigma, \\ x^*, y^* \in \partial \sigma^* \quad (37)$$

Solving (37) with respect to δu_n and substituting the result in (36), obtain

$$\delta J_\psi = \int_0^T \int_{\partial \sigma} \{ [\psi' f_{u_x} - (\psi' f_{u_x})_x] \cos \nu + [\psi' f_{u_y} - (\psi' f_{u_y})_y] \sin \nu + \frac{\partial}{\partial S} [\psi' f_{u_x} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) g^{-1} u_n g_u] \delta u_n ds dt + \int_0^T \int_{\partial \sigma} \cos \nu - [(\psi' f_{u_x} - (\psi' f_{u_x})_x) \cos \nu + (\psi' f_{u_y} - (\psi' f_{u_y})_y) \sin \nu] u_x - \frac{\partial}{\partial S} (\psi' f_{u_x} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu) u_x - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) g^{-1} u_n g_x - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) \left(\frac{\partial^2 u}{\partial n^2} \cos \nu - \frac{\partial^2 u}{\partial n \partial S} \sin \nu \right) \} \delta x ds dt + \int_0^T \int_{\partial \sigma} \{ S(x, y, t) \sin \nu - [(\psi' f_{u_x} - (\psi' f_{u_x})_x) \cos \nu + (\psi' f_{u_y} - (\psi' f_{u_y})_y) \sin \nu] u_y - \frac{\partial}{\partial S} (\psi' f_{u_x} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu) u_y - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) g^{-1} u_n g_y - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) \left(\frac{\partial^2 u}{\partial n^2} \sin \nu + \frac{\partial^2 u}{\partial n \partial S} \cos \nu \right) \} \delta y ds dt, (x, y) \in \partial \sigma \quad (38)$$

We now choose boundary conditions for equation (30), so that the first term on the right-hand side of expression (38) vanishes. In view of the arbitrariness of the variation δu on the contour, obtain

$$[\psi' f_{u_x} - (\psi' f_{u_x})_x] \cos \nu + [\psi' f_{u_y} - (\psi' f_{u_y})_y] \sin \nu + \frac{\partial}{\partial S} [\psi' f_{u_x} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu] - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) g^{-1} u_n g_u = 0. \quad (39)$$

Substituting (30) into (29) and changing the order of integration, obtain the final expression

$$\delta J_\psi = \int_{\partial \sigma} L_1(x, y) \delta x(x, y) ds + \int_{\partial \sigma} L_2(x, y) \delta y(x, y) ds, \quad (40)$$

where

$$L_1(x, y) = \int_0^T \{ S(x, y, t) \cos \nu - [(\psi' f_{u_x} - (\psi' f_{u_x})_x) \cos \nu + (\psi' f_{u_y} - (\psi' f_{u_y})_y) \sin \nu] u_x - \frac{\partial}{\partial S} (\psi' f_{u_x} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu) u_x - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) g^{-1} u_n g_x - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) \left(\frac{\partial^2 u}{\partial n^2} \cos \nu - \frac{\partial^2 u}{\partial n \partial S} \sin \nu \right) \} dt, \quad (41)$$

$$L_2(x, y) = \int_0^T \{ S(x, y, t) \sin \nu - [(\psi' f_{u_x} - (\psi' f_{u_x})_x) \cos \nu + (\psi' f_{u_y} - (\psi' f_{u_y})_y) \sin \nu] u_y - \frac{\partial}{\partial S} (\psi' f_{u_x} \sin \nu \cos \nu - \psi' f_{u_{yy}} \sin \nu \cos \nu) u_y - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) g^{-1} u_n g_y - (\psi' f_{u_{xx}} \cos^2 \nu + \psi' f_{u_{yy}} \sin^2 \nu) \left(\frac{\partial^2 u}{\partial n^2} \sin \nu + \frac{\partial^2 u}{\partial n \partial S} \cos \nu \right) \} dt, (x, y) \in \partial \sigma \quad (42)$$

Note that the perturbation of the contour $\delta(\partial \sigma)$ is represented by perturbation of the corresponding coordinates δx and δy contour at each point of the originally defined boundary. The necessary conditions for optimality (of the first order) will be the conditions:

$$L_1(x, y) = 0, \quad (43)$$

$$L_2(x, y) = 0. \quad (44)$$

IV. CONCLUSIONS

The article discovered necessary optimality conditions for determining of the position of the boundary of oil deposit. The essential criteria was discovered in the thesis to define free boundary over the nonlinear partial differential equations.

REFERENCES

- [1] Zakirov SN, Lapuk BB. Design and development of gas fields. Moscow: Nedra, 1974.
- [2] Lions J-L. Optimal control of systems described by partial differential equations. Moscow: Mir, 1972.
- [3] Bubnov BA. The inverse problem for a parabolic equation and the Cauchy problem for certain classes of evolution equations in classes of functions of finite smoothness. Dokl. AN SSSR, Moscow. 1988;299(5):782-784.
- [4] Bulygin VYa. Hydrodynamic analysis of the exploited oil field by solving the correct and ill-posed problems. In: Theory and experiment: problems of development

- of the field. Kazan: Publishing House of Kazan University, 1972.
- [5] Bulygin VYa. Hydromechanics of the oil reservoir. Moscow: Nedra, 1974.
- [6] Gutnikov AI, Zholdasov A, Zakirov SN. Deposits of gas and oil in natural and artificial streams of water. All-Russian Research Institute of Gas Industry, Seria: Geol. and gas condensate fields, Moscow. 1986;4:54.
- [7] Gutnikov AI, Zholdasov A, Zakirov SN. The configurations of gas and oil deposits in hydrodynamic traps. Higher Education Messenger: Oil and Gas, Moscow. 1984;11:3-6.
- [8] Gutnikov AI, Zholdasov A, Zakirov SN. Nonstationary displacements of an oil deposit in a stream of stratal water. Messenger of Academy of Sciences of the USSR, Moscow. 1985;2:177-179.
- [9] Zhiron VV. New methods for forecasting the composition of extracted production in the development of natural gas deposits. PhD Thesis. Moscow: VNII Gas, 1985.
- [10] Shazhdekeeva NK, Mukhambetzhano ST. On the properties of the solution of a certain problem of the theory of filtration with respect to phase transitions. Materials of the V international conference "Mathematical modeling and information technologies in education and science", dedicated to the 25th anniversary of computer science in the school. Almaty, 2010, pp. 114-118.
- [11] Shazhdekeeva NK, Kammatov K, Adieva AA. On the existence of stationary solutions of a certain type of quasilinear systems. Messenger of ASU, Atyrau. 2005;2:31-36.
- [12] Frolov VV. Uniqueness theorems for the solution of the inverse heat conduction problem. Physics Engineering Journal. 1975;29(1):145-150.